

## Time-resolved spectroscopy of the planet-hosting sdB pulsator V391 Pegasi

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### Abstract

The subdwarf B (sdB) star V391 Peg oscillates in short-period p modes and long-period g modes, making it one of the three known hybrids among sdBs. As a by-product of the effort to measure secular period changes in the p modes due to evolutionary effects on a time scale of almost a decade, the O–C diagram has revealed an additional sinusoidal component attributed to a periodic shift in the light travel time caused by a planetary-mass companion around the sdB star in a 3.2 yr orbit. In order to derive the mass of the companion object, it is necessary to determine the orbital inclination. One promising possibility to do this is to use the stellar inclination as a primer for the orbital orientation. The stellar inclination can refer to the rotational or the pulsational axis, which are assumed to be aligned, and can in turn then be derived by combining measurements of  $v_{rot}$  and  $v_{rot} \sin i$ .

The former is in principle accessible through rotational splitting in the photometric frequency spectrum (which has however not been found for V391 Peg yet), while the projected rotational velocity can be measured from the rotational broadening of spectral lines. The latter must be deconvolved from the additional pulsational broadening caused by the surface radial velocity variation in high S/N phase averaged spectra. This work gives limits on pulsational radial velocities from a series of phase resolved spectra.

Phase averaged and phase resolved high resolution echelle spectra were obtained in May and September 2007 with the 9m-class Hobby-Eberly Telescope (HET), and one phase averaged spectrum in May 2008 with the 10m-Keck 1 telescope<sup>1</sup>.

Individual Objects: V391 Pegasi

### The hybrid pulsating sdB star V391 Pegasi and its planetary companion

Subdwarf B stars (sdBs) are subluminescent, evolved stars on the extreme horizontal branch (EHB). They have a He burning core but, due to previous significant mass loss, no H-shell burning in their thin hydrogen shells. Their masses cluster around  $0.5 M_{\odot}$ . Only a small fraction of the sdBs show pulsational variations, with non-pulsators also populating the region in the HRD where the pulsators are found. There are p (pressure) mode and g (gravity) mode types of pulsation. Three objects that show both mode types are referred to as hybrid pulsators and among them is V391 Peg (HS 2201+2610) which has five p modes (Østensen et al. 2001; Silvotti et al. 2002) and one g mode (Lutz et al. 2008, 2009). Silvotti

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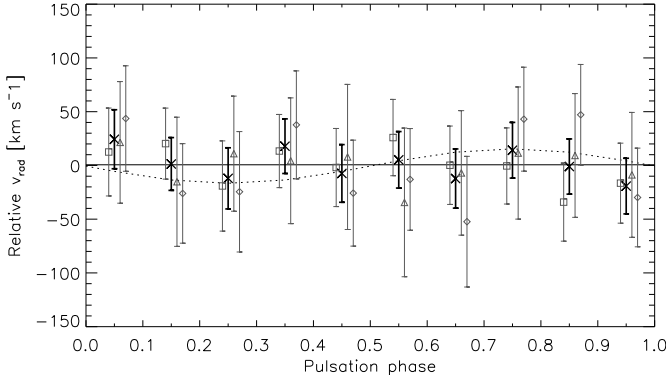


Figure 1: Radial velocities derived for the  $H\beta$  (squares),  $H\gamma$  and  $H\delta$  (triangles and diamonds) lines from cross correlation with the template (with error bars, and plotted at a small phase offset for clarity); and weighted mean of the results (crosses and thick error bars) together with a constant and a sinusoidal curve.

et al. (2007) detected parabolic and sinusoidal variations in the observed–calculated (O–C) diagram constructed for the two main pulsation frequencies at 349.5 s and 354.2 s over the observing period of seven years. The sinusoidal component is attributed to the presence of a very low-mass companion (V391 Peg b,  $m \sin i = 3.2 \pm 0.7 M_{\text{Jup}}$ ). The determination of the true mass of this ‘asteroseismic planet’ requires a constraint on the orbital inclination which can presumably be determined via the stellar rotational inclination.

### Limits on the pulsational radial velocities from phase resolved echelle spectra

Echelle spectra of V391 Peg were taken during May and September 2007 with the HRS ( $R = 15000$ ) of the HET at the McDonald Observatory, and with HIRES ( $R = 31000$ ) at the Keck 1 telescope atop Mauna Kea in May 2008. Data reduction was done with ESO-Midas using standard procedures. Individual echelle orders were merged and the final spectra were carefully normalized and finally summed (Kruspe 2009, diploma thesis, in prep.). This results in a set of individual spectra ( $S/N \approx 3$ ), in particular two September 2007 high time resolution series, and summed spectra for May and September 2007.

In our attempt to “clean” the relevant rotational broadening from pulsational effects, the spectra in September obtained in time resolved mode were combined to a series of ten phase resolved averaged spectra ( $S/N \approx 9$ ) for the main pulsation period of 349.5 s (similar to Tillich et al. 2007).

The cross-correlation of this series of averaged spectra with a pure hydrogen NLTE model spectrum at  $T_{\text{eff}} = 30000$  K and  $\log(g/\text{cm s}^{-2}) = 5.5$  as a template yields pulsational radial velocity measurements as shown in Fig. 1 for the different Balmer lines. The maximum amplitude of a sinusoidal curve (fixed at the expected period) that could be accommodated in comparison to the weighted means of the Balmer lines reveals that any pulsational radial velocity amplitude is smaller than the accuracy of our measurements and confirms the upper limit of  $16 \text{ km s}^{-1}$  given by Kruspe et al. (2008).

The resolution of the model template matches the spectral resolution of the (pulsation-averaged) Keck spectrum. A comparison of the  $H\alpha$  NLTE line core shape yields an even more stringent upper limit for the combined broadening effect of pulsation and rotation of at most  $9 \text{ km s}^{-1}$ . This means that much better data in terms of spectral resolution and signal to noise will be necessary to measure  $v_{\text{puls}}$  and  $v_{\text{rot}} \sin i$ .

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